# Jet medium modifications

#### Carlota Andres

#### CPHT, École polytechnique HP2023, Aschaffenburg, March 26 -31







#### Jet modification in HI

- Medium-induced energy loss
  - Out-of-cone energy loss
  - Jet and hadron suppression





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• During the formation time of the gluon **multiple scatterings** act **coherently** 



Landau, Pomeranchuk, Migdal

Suppression of the spectrum for large formation times

• Resummation of multiple scatterings: **BDMPS-Z formalism** 

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  - For energy loss calculations, we only need the soft limit ( $z \ll 1$ ):

Formally :  $E \to \infty, z \to 0$  ( $\omega = zE$  finite) Emission kernel Broadening  $\omega \frac{dI}{d\omega d^2 k} = \frac{2\alpha_s C_R}{(2\pi)^2 \omega^2} \operatorname{Re} \int_0^\infty dt' \int_0^{t'} dt \int_{pq} p \cdot q \underbrace{k(t', q; t, p)}_{\text{Baier, Dokshitzer, Mueller, Peigné, Schiff (96)}}_{\text{Zakharov (97)}}$ 

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• In practice, solved for some approximations

• **Opacity expansion** in the number of scatterings

N = 1: GLV Gyulassy, Levai, Vitev (2000)

- Harmonic oscillator (HO): multiple <u>soft</u> scatterings (Gaussian approximation)  $n\sigma(r) \approx \frac{1}{2}\hat{q}r^2$
- AMY: <u>infinite length</u> medium Arnold, Moore, Yaffe (2002)
- Recent approaches going beyond these approximations
  - Improved opacity expansion (IOE) Semi-analytical expansion around the HO
  - Fully resummed spectrum ( $z \ll 1$ )

Kernel as a time dependent Schrödinger equation

tions  $V(q) \propto \frac{1}{q^4}$ Mehtar-Tani, Barata, Soto-Ontoso, Tywoniuk,

<u>1903.00506</u>, <u>2004.02323</u>, <u>2106.07402</u>

 $i\sigma(r)$ 

CA, Apolinario, Martinez, Dominguez, <u>2002.01517</u>, <u>2011.06522</u>

Beyond the soft limit but integrated in  $k_T$ : Caron-Huot and Gale, <u>1006.2379</u>





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### Medium-induced radiation: nonperturbative potential

- Perturbative potential (broadening kernel) known at LO and NLO in pQCD Aurenche Gelis Zaraket (2002), Caron -Huot (2008)
- Non-perturbative result from EQCD Moore, Schlusser, <u>2009.06614</u> Ghiglieri's talk
- Matching between the EQCD IR and the perturbative UV results Moore, Schlichting, Schlusser, Soudi, <u>2105.01679</u>
   Finite length rates



(z)E

E

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30
12
60

Yazdi, Shi, Jeon, Gale, <u>2206.05855</u>

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#### From energy loss to jet substructure

• Grooming techniques to isolate prongs corresponding to a hard splitting



• Soft limit ( $z \ll 1$ ) not that useful:

 $\frac{\mathrm{dI}^{\mathrm{med}}}{\mathrm{d}z\mathrm{d}\theta}??$ 

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#### Beyond the soft limit



• Complete (multiple scatterings) medium-induced emission spectrum keeping z and  $\theta$  not available

until last Thursday! (See Isaksen's talk Tuesday 14:40) and slide 12 of this talk

- Two available approaches:
  - Opacity expansion:
    - *N* = 1 result Ovanesyan, Vitev <u>1103.1074</u>, <u>1109.5619</u>
    - Highly complicated recursive relations to go to all orders

Sievert, Vitev, <u>1807.03799</u>

See also

Djordjevic's talk

Tues 09:20

- Some numerical results at N = 1 and N = 2 Sievert, Vitev, Yoon, <u>1903.06170</u>
- Unitarity issues can lead to negative cross sections
- *Tilted* Wilson lines:
  - Assumes <u>semi-hard</u> splittings (z not too small)
  - All partons propagate along straight line trajectories

Dominguez, Milhano, Salgado, Tywoniuk, Vila, <u>1907.03653</u>

Isaksen, Tywoniuk 2107.02542



• **Two-point energy correlator** of a heavy-ion quark jet:

Proposed by CA, Dominguez, Elayavalli, Holguin, Marquet, Moult

$$\frac{\mathrm{d}\Sigma^{(n)}}{\mathrm{d}\theta} = \frac{1}{\sigma_{qg}} \int \mathrm{d}z \, \frac{\mathrm{d}\sigma_{qg}}{\mathrm{d}z\mathrm{d}\theta} \, z^n (1-z)^n + \mathcal{O}\left(\frac{\mu_s}{E}\right)$$



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$$\frac{\mathrm{d}\Sigma^{(n)}}{\mathrm{d}\theta} = \left(\frac{1}{\sigma_{qg}}\int \mathrm{d}z \left(\frac{\mathrm{d}\sigma_{qg}^{\mathrm{vac}}}{\mathrm{d}z\mathrm{d}\theta} + \frac{\mathrm{d}\sigma_{qg}^{\mathrm{med}}}{\mathrm{d}z\mathrm{d}\theta}\right) z^{n}(1-z)^{n}\right) \left(1 + \mathcal{O}\left(\frac{\bar{\mu}_{s}}{Q}\right)\right) + \mathcal{O}\left(\frac{\Lambda_{\mathrm{QCD}}}{\theta Q}\right)$$
$$\frac{1}{\sigma_{qg}}\frac{\mathrm{d}\sigma_{qg}^{\mathrm{vac}}}{\mathrm{d}z\mathrm{d}\theta} = \frac{\alpha_{s}(\theta Q)}{\pi}C_{F}\frac{1 + (1-z)^{2}}{z\theta} + \mathcal{O}\left(\alpha_{s}^{2},\theta\right)$$

Inclusive medium-induced cross section differential in z and  $\theta$ 



• **Two-point energy correlator** of a heavy-ion quark jet:

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Also experimental talks: <u>Cruz-Torres</u> for ALICE Tues 17:30, and <u>Tamis</u> for STAR Wed 11:30

### Beyond the soft limit: $g \rightarrow c\bar{c}$

• Not divergent for  $z \to 0$ . Need to go beyond the soft limit

Results with multiple scatterings but integrated in transverse momentum: Caron-Huot and Gale, <u>1006.2379</u>

• Keeping both z and  $\theta$ :



• Opacity expansion results at N = 1 and N = 2 Sievert, Vitev, Yoon, <u>1903.06170</u>

• Results with **resummation of multiple scatterings** recently obtained in the large  $N_c$  limit and HO approximation

(No need of the semi-hard approximation)



### Beyond the soft limit: $\gamma \rightarrow q\bar{q}$

- Results with multiple scatterings in the semi-hard approximation available at large and finite N<sub>c</sub>
   Dominguez, Milhano, Salgado, Tywoniuk, Vila, <u>1907.03653</u> Isaksen, Tywoniuk <u>2107.02542</u>
- **Complete** (multiple scatterings) medium-induced emission spectrum **keeping** *z*

and  $\theta$  (Without using of the *semi-hard* approximation!!) Isaksen, Tywoniuk, <u>2303.12119</u> (HO approximation)



Isaksen, Tywoniuk, <u>2303.12119</u>

 $\boldsymbol{E}$ 

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-z)E

### Medium-induced radiation and transverse dynamics

• Jets decouple from the medium transverse dynamics in the usual (*eikonal*) medium-induced approaches



- Need of **generalizing** the medium-induced formalisms to account for  $\mathcal{O}(1/E)$  (*subeikonal*) terms
- GLV emission spectrum (and broadening) obtained: Sadofyev, Sievert, Vitev, 2104.09513
  - for transversely **flowing** matter
  - for transversely **inhomogeneous** matter

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#### Medium-induced radiation and transverse dynamics

- Broadening with multiple scatterings obtained:
  - for transversely flowing matter CA, Dominguez, Sadofyev, Salgado, <u>2207.07141</u>
  - for transversely **inhomogeneous** matter

See also Barata, Sadofyev, Wang 2210.06519

#### Transverse inhomogeneities

Fu, Casalderrey-Solana, Wang, 2204.05323

Barata, Sadofyev, Salgado, <u>2202.08847</u> Fu, Casalderrey-Solana, Wang, <u>2204.05323</u> See also Mayo López's poster

Transverse flow (single scattering)





 $\omega = 5 \text{ GeV}$ 1.0 )=10 GeV ω=20 GeV 0.5 -<p<sup>3</sup>> (GeV<sup>3</sup>) 0.0 -0.5 -1.0 -15 -10 -5 0 5 10 15 Initial transverse position x (fm) Xin-Nian Wang's talk Tue 10:00

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Sadofyev's talk Wed 10:00

• Overlapping effects among multiple splittings?

Extremely hard calculation performed for  $g \rightarrow ggg$ (on-shell emitter, infinite, static medium, large  $N_c$ , HO)



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Seems to be a small effect for gluon showers. Parton showers seem to be safe!

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Arnold, Iqbal, Chang, Gorda, Rase, Elgeadwy (2015-2022) 2212.08086, 2302.10215

#### • Independent medium-induced splittings

Jeon, Moore (2003) Blaizot, Dominguez, Iancu, Mehtar-Tani (2014)

$$\frac{\partial}{\partial t}D(x,t) = \frac{1}{t_*}\int \mathrm{d}z\,\mathcal{K}(z)\left[\sqrt{\frac{z}{x}}D\left(\frac{x}{z},t\right) - \frac{z}{\sqrt{x}}D\left(x,t\right)\right]_{\text{BDIM}}$$

$$\frac{\partial}{\partial t}D(x,\boldsymbol{k},t) = \frac{1}{t_*}\int \mathrm{d}z\,\mathcal{K}(z)\left[\frac{1}{z^2}\sqrt{\frac{z}{x}}D\left(\frac{x}{z},\frac{\boldsymbol{k}}{z},t\right) - \frac{z}{\sqrt{x}}D\left(x,\boldsymbol{k},t\right)\right] + \int_{\boldsymbol{q}}\mathcal{C}(\boldsymbol{q})D(x,\boldsymbol{k}-\boldsymbol{q},t)$$

\*Neglecting any momentum transfer during the formation time of the splittings

 $t_* = \alpha \sqrt{\frac{E}{\hat{a}}}$ 

• Overlapping effects among multiple splittings?

Extremely hard calculation performed for  $g \rightarrow ggg$ (on-shell emitter, infinite, static medium, large  $N_c$ , HO)



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Arnold, Iqbal, Chang, Gorda, Rase, Elgeadwy (2015-2022) 2212.08086, 2302.10215

• Independent medium-induced splittings  
Jeon, Moore (2003) Blaizot, Dominguez, Iancu, Mehtar-Tani (2014)  

$$\frac{\partial}{\partial t} D(x,t) = \frac{1}{t_*} \int dz \, \mathcal{K}(z) \left[ \sqrt{\frac{z}{x}} D\left(\frac{x}{z},t\right) - \frac{z}{\sqrt{x}} D\left(x,t\right) \right] \text{BDIM} t_* = \alpha \sqrt{\frac{E}{\hat{q}}}$$

$$\frac{\partial}{\partial t} D(x,k,t) = \frac{1}{t_*} \int dz \, \mathcal{K}(z) \left[ \frac{1}{z^2} \sqrt{\frac{z}{x}} D\left(\frac{x}{z},\frac{k}{z},t\right) - \frac{z}{\sqrt{x}} D\left(x,k,t\right) \right] + \int_q^{C} C(q) D(x,k-q,t)$$

\*Neglecting any momentum transfer during the formation time of the splittings



• Generalization of BDMI to account for the transmission of **polarization** in an **anisotropic** medium (Static medium, HO,  $\hat{q}$  constant but  $\hat{q}_y \neq \hat{q}_z$ )

Polarization is found to be created and washed out at each step of the branching Hauksson, Iancu 2303.03914

- Energy and angular profile of the medium-induced cascade
  - Kinetic theory approach: linearized Boltzmann equation
     Energy distribution: Soudi, Schlichting 2008.04928
     Angular distribution: Mehtar-Tani, Soudi, Schlichting, 2209.10569
  - Solving **BDMI** equations for **longitudinally expanding media** (HO)

Energy distribution: Adhya, Spousta, Salgado, Tywoniuk, <a href="https://www.2106.02592"><u>Adhya's poster</u></a>Angular distribution: Adhya, Kutak, laczek, Rohrmoser, Tywoniuk, <a href="https://www.2211.15803">2211.15803</a>





**Out-of-cone** energy loss via **medium-induced radiation**, followed by **elastic scatterings** of soft fragments pushing the distribution to **large angles** and thermalization

### Jet quenching in the initial stages?

- Jet quenching is the only QGP signature not (yet?) observed in small systems
- In small systems the **initial stages** (IS) are specially important

#### Understanding jet quenching in the initial stages becomes crucial!

- Usually, jet quenching analyses neglect energy loss in the IS, but some observables have been found sensitive to the IS
   CA, Armesto, Niemi, Paatelainen, Salgado, <u>arXiv:1902.03231</u>
   Stojku, Auvinen, Djordjevic, Huovinenm, Djordjevic, <u>arXiv:2008.08987</u>
- Many new developments in the computation of the broadening in the pre-hydro stages
  - In the Glasma phase: Ipp, Müller, Schuh, <u>arXiv:2001.10001</u>, <u>arXiv:2009.14206</u>

Carrington, Czajka, Mrówczynski, arXiv:2112.06812, arXiv:2202.00357

Avramescu, Baran, Greco, Ipp, Müller, Ruggieri, 2303.05599 Avramescu's talk

Tues 12:10

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• Between the Glasma and hydro phases with Kinetic theory

Boguslavski, Kurkela, Lappi, Lindenbauer, Peuron, <u>2303.12595</u>

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#### Understanding jet quenching in the initial stages becomes crucial!



#### $\hat{q}$ relatively large!

#### Summary

- Many new developments in the theory of **medium-induced radiation** 
  - Improved analytical approximations and numerical solutions for the emission spectrum in the soft limit ( $z \ll 1$ )
  - Non-perturbative determination of the potential (broadening kernel)
  - Results **beyond the soft limit keeping the angle**:
    - Opacity expansion
    - Semi-hard approximation for multiple scatterings
    - First complete results with multiple scatterings
  - Calculations coupling the jet to the **medium's transverse dynamics**
- Overlapping effects between multiple emissions in gluon showers found to be small (with caveats)
- Calculations of the **broadening** in the **pre-hydrodynamics** stages

# Thank you!

• The emitted gluon carries a fraction of energy *z* of the energy of the parent quark

$$\begin{aligned} z \frac{dI}{dz d^2 \mathbf{k}} &= \frac{\alpha_s P_{g \leftarrow q}(z)}{(2\pi)^2 z (1-z)^2 E^2} \operatorname{Re} \int_0^\infty dt' \int_0^{t'} dt \int_{\mathbf{p}_1 \mathbf{p}_2 \mathbf{k}_1 \mathbf{k}_2} \mathcal{P}_q(t, \mathbf{p}_1; 0, \mathbf{p}_0) \\ &\times (\mathbf{k}_1 - z \mathbf{p}_1) \cdot (\mathbf{k}_2 - z \mathbf{p}_2) \ \widetilde{\mathcal{K}}^{(3)}(t', \mathbf{k}_2 - z \mathbf{p}_2; t, \mathbf{k}_1 - z \mathbf{p}_1; \mathbf{p}_2 - \mathbf{p}_1) \\ &\times \mathcal{P}_g(\infty, \mathbf{k}; t', \mathbf{k}_2) \end{aligned}$$

Blaizot, Dominguez, Iancu, Mehtar-Tani arXiv:1209.4585 Apolinario, Armesto, Milhano, Salgado arXiv:1407.0599



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• The emitted gluon carries a fraction of energy *z* of the energy of the parent quark

$$z \frac{dI}{dz d^2 k} = \frac{\alpha_s P_{g \leftarrow q}(z)}{(2\pi)^2 z (1-z)^2 E^2} \operatorname{Re} \int_0^\infty dt' \int_0^{t'} dt \int_{p_1 p_2 k_1 k_2} \mathcal{P}_q(t, p_1; 0, p_0)$$
  
  $\times (k_1 - z p_1) \cdot (k_2 - z p_2) \widetilde{\mathcal{K}}^{(3)}(t', k_2 - z p_2; t, k_1 - z p_1; p_2 - p_1)$   
  $\times \mathcal{P}_g(\infty, k; t', k_2)$ 

Blaizot, Dominguez, Iancu, Mehtar-Tani arXiv:1209.4585 Apolinario, Armesto, Milhano, Salgado arXiv:1407.0599



• The emitted gluon carries a fraction of energy *z* of the energy of the parent quark

$$z\frac{dI}{dzd^{2}\boldsymbol{k}} = \frac{\alpha_{s}P_{g\leftarrow q}(z)}{(2\pi)^{2}z(1-z)^{2}E^{2}}\operatorname{Re}\int_{0}^{\infty} dt'\int_{0}^{t'} dt\int_{\boldsymbol{p}_{1}\boldsymbol{p}_{2}\boldsymbol{k}_{1}\boldsymbol{k}_{2}} \mathcal{P}_{q}(t,\boldsymbol{p}_{1};0,\boldsymbol{p}_{0})$$

$$\times (\boldsymbol{k}_{1}-z\boldsymbol{p}_{1}) \cdot (\boldsymbol{k}_{2}-z\boldsymbol{p}_{2})\widetilde{\mathcal{K}}^{(3)}(t',\boldsymbol{k}_{2}-z\boldsymbol{p}_{2};t,\boldsymbol{k}_{1}-z\boldsymbol{p}_{1};\boldsymbol{p}_{2}-\boldsymbol{p}_{1})$$

$$\times \mathcal{P}_{g}(\infty,\boldsymbol{k};t',\boldsymbol{k}_{2})$$
Blaizot, Dominguez, Iancu, Mehtar-Tani arXiv:1209.4586  
Apolinario, Armesto, Milhano, Salgado arXiv:1407.0599
$$p_{0}, \mathbf{x}' \qquad p_{1} \qquad p_{1} - \mathbf{k}_{1} - z \qquad \mathbf{k}'$$

$$p_{0}, \mathbf{x}' \qquad p_{2} \qquad \mathbf{k}'$$

$$\mathbf{I} \qquad \mathbf{I} \qquad$$

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$$\times (\boldsymbol{k}_{1} - z\boldsymbol{p}_{1}) \cdot (\boldsymbol{k}_{2} - z\boldsymbol{p}_{2}) \widetilde{\mathcal{K}}^{(3)}(t', \boldsymbol{k}_{2} - z\boldsymbol{p}_{2}; t, \boldsymbol{k}_{1} - z\boldsymbol{p}_{1}; \boldsymbol{p}_{2} - \boldsymbol{p}_{1})$$

$$\times \mathcal{P}_{g}(\infty, \boldsymbol{k}; t', \boldsymbol{k}_{2})$$
Blaizot, Dominguez, Iancu, Mehtar-Tani arXiv:1209.4588  
Apolinario, Armesto, Milhano, Salgado arXiv:1407.0599
$$p_{0}, \mathbf{x}' \qquad p_{1} \qquad p_{1} - \mathbf{k}_{1}, 1 - z \qquad \mathbf{k}'_{2} \qquad \mathbf{k}'_{1} \qquad \mathbf{k}'_{2} \qquad \mathbf{k}'_{1} \qquad \mathbf{k}'_{2} \qquad \mathbf{k}'_{1} \qquad \mathbf{k}'_{1} \qquad \mathbf{k}'_{2} \qquad \mathbf{k}'_{1} \qquad \mathbf{k}'_{2} \qquad \mathbf{k}'_{1} \qquad \mathbf{k}'_{2} \qquad \mathbf{k}'_{1} \qquad \mathbf{k}'_{2} \qquad \mathbf{k}'_{$$

#### Non-perturbative potential



Yazdi, Shi, Jeon, Gale, <u>2206.05855</u>

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