Slow modes and momentum anisotropy in far-from-equilibrium QGP

Introduction

- Far-from-equilibrium slow modes from kinetic theory and their relations to hydro. modes.
 - Novel slow modes associated with momentum anisotropy.
- Implications for the origin of flow in small systems?

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Weiyao Ke @ LANL

Brewer@ CERN

Far-from-equilibrium QGP

- Heavy-ion collision: a unique opportunity for studying offequilibrium properties of quark-gluon matter.
- Significant progress in theory (attractor, KOMPOST etc) but important questions still open: e.g. what are the relevant d.o.fs at early time.

Works by many; Review: Burgers et al, Rev.Mod.Phys. 2020

- Near equilibrium: slow modes are conserved densities as described by hydro. → the evolution of azimuthal anisotropy is driven by hydro. modes.
- Far from equilibrium: are there slow modes? how does azimuthal anisotropy evolve (the origin of flow in small systems)? This talk: addressing the questions within kinetic theory models.

Multipole Moments

• Consider the phase space distribution $f(\vec{p}; \tau)$ of gluons in an Bjorken expanding plasma

$$L_{l,m}(\tau) \equiv \int_{\vec{p}} p P_l^m(\theta) \cos(m\phi) f(\vec{p};\tau)$$

e.g. $L_{0,0} = T^{00}, L_{2,0} \propto \epsilon - p_L/3, L_{2,1} = T^{0x}, L_{2,2} \propto T^{xx} - T^{yy}.$

- m: the shape of the distribution in transverse plane ($m \ge 2$ measures the azimuthal anisotropy)
- *l*: the correlation strength at scale $\Delta v_z \sim \cos(\pi/l)$.
- Collecting $L_{l,m}$ with the same m and parity $(z \rightarrow -z) s = \pm$ into "a bigger vector" $\Psi_{m,\pm}$,

• e.g
$$\psi_{0+} = (L_{0,0}, L_{2,0}, \dots), \psi_{2+3} = (L_{2,2}, L_{4,2}, \dots)$$

RTA Kinetic theory

• Consider the angular distribution $F(\tau, \hat{p}) = \int dp p^2 p f(\tau, \vec{p})$ and its evolution eqn

$$\partial_{\tau}F - \frac{\hat{p}_z(1-\hat{p}_z^2)}{\tau}\partial_{\hat{p}_z}F + \frac{4\hat{p}_z^2}{\tau}F = -C[f]$$

 Using the generalized relaxation time approximation collision integral (conserving energy-momentum) :

m

$$C[F] = -\frac{F - T^{00} - 3(T^{0x}\hat{p}_x + T^{0y}\hat{p}_y + \tau T^{0\eta}\hat{p}_z)}{\tau_R},$$

ne evolution eqns for moments (different sectors (ms) do not ix).

 $\partial_{\tau} \psi_{ms} = \mathcal{H}_{ms}(\tau) \psi_{ms}$

Slow modes

• The collection of low-lying eigenmodes which are gapped from the others.

$$\mathscr{H}(\tau)\phi(\tau) = E(\tau)\phi(\tau)$$

non-Hermitian complex



(decay rate $ReE \ge 0$ because of expansion and collision)

• Dominate the dynamics under adiabatic conditions.

Brewer-Li Yan-YY, PLB 21', Brewer-Scheihing-Hitschfeld-YY, JHEP 22', Mikheev-Mazeliauskas-Berges, PRD 22';



• There is an early-time slow mode associated with energy density.

$$\phi_{0+}^G = (T^{00}, L_{20}, \dots) \qquad \phi_{0+}^G = (T^{00}, 0, \dots)$$

Similar results for (I+) sector (containing T^{0x}); one-to-one correspondence between conserved density and early-time slow modes?



- Early-time slow mode exists for $m \ge 2$ (anisotropy) sectors.
- Democracy among different m(s) at early time:

$$\lim_{\tau \to 0} E_{0+}^G = E_{1+}^G = E_{2+}^G = E_{3+}^G = \dots = \frac{1}{\tau}$$

Shapes in phase space as slow modes



- The phase space volume is preserved (Liouville theorem) in collisionless limit. (c.f. 2203.05004).
 - For Bjorken expansion at early time, the shape of transverse distribution evolves slowly.

Nonlinear Bosonization of Fermi Surfaces: The Method of Coadjoint Orbits

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Collisionless dynamics of general non-Fermi liquids from hydrodynamics of emergent conserved quantities

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Many recent developments on emergent symmetries (slow modes), dynamics of shape in collisionless kinetic theory.

Slow modes in and out of equilibrium



far-from-equilibrium (large $1/\tau$)

near-equilibrium (small $1/\tau$)

- Showing the diverse relation to conserved densities.
- Implications for the memory of initial azimuthal an-isotropy?

The evolution rate

$$V_m^{(n)} = \int_{\hat{p}} (\frac{p_T}{p})^n F(\hat{p}) \cos(m\phi)$$

Harmonics
Given the dominance of slow modes,
the percentage change rate of $V_m^{(n)}$ is
insensitive to many details of initial
conditions (attractor behavior).
$$\tau \partial_{\tau} \log V_m^{(n)} = -E_{m+}^G$$

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Ground state
eigenvalue

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 $-1.6 \frac{1}{10^{-2}}$

rate of V_m

 10^{-1}

 τ/τ_R

m = 2

 10^{0}

Initial azimuthal anisotropy decay slowly m = 2

0.5

0.4

0.2

Ξ_{E 0.3}

$$u_m^{(n)} = \frac{V_m^{(n)}}{V_0^{(n)}} \, (m \ge 2)$$

NB: u_2, u_3 is the proxy of elliptic, triangle flow etc. $u_2^{(2)} \propto (T^{xx} - T^{yy})/\epsilon$.



- The memory of initial momentum space eccentricity will survive up to τ_R
 - Without $m \ge 2$ slow modes, u_m would quickly decay as $(\tau/\tau_I)^{-\#}$.

Summary and outlook

- We study the slow modes of the far-from-equilibrium QGP using kinetic theory.
 - The relation between slow modes and conserved density is rather diverse.
 - Momentum space an-isotropy (the shape of distribution) evolve slowly.
- Implication: accounting for the prehydro. evolution of azimuthal an-isotropy in terms of slow modes. (Further check is needed, e.g. including spatial gradient effects, more general microscopic model).

A broader view: the evolution of QGP vs scale



- This talk: QGP properties from equilibrium to far-fromequilibrium (QGP with varying expansion rate).Other perspectives:
 - Jet observables as a function of scattering angle.
 - Medium response with varying gradient (Jet-medium responses).

Ultimate goal: a comprehensive picture of QGP evolution!

Back-up

Attractor



- After the decay of fast modes, the ratio of moment is determined by that of slow mode and is insensitive to the initial condition (attractor behavior).
- Early-time attractor behavior for higher spin sectors.
- Implications for the memory of momentum space anisotropy?

Longitudinal momentum density sector

 $(ms)=(0-), e.g. L_{01} = T^{0z}$



• T^{0z} is not related to the early-time slow modes.